Subwavelength Grating Filters in SOI

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Abstract: We demonstrate subwavelength grating filters in silicon-on-insulator based on Bragg gratings and racetrack resonators. The Bragg grating has a 3 dB bandwidth = 0.5 nm and reflectivity > 90%; the resonator has a 3 dB bandwidth = 1 nm and extinction ratio > 33 dB.

1. Introduction

There is an increasing need for integrated solutions in optical communications, interconnections, and sensing applications [1]. In the past few years, a variety of active and passive devices as well as integrated subsystems in CMOS-compatible silicon photonics platforms have been realized [1,2]. One specific structure that has attracted considerable interest is the subwavelength grating (SWG) [3]. SWG waveguides have the potential for low loss and provide significant flexibility in tailoring the effective index. Indeed, a number of SWG-based building blocks in silicon-on-insulator (SOI) have been developed, including waveguide crossings, bends, couplers, mode transformers, and polarization converters [4]. Recently, we proposed the designs of SWG filters based on Bragg gratings (BGs), microring resonators, and racetrack resonators [5]; however, to date, only microring resonators have been verified experimentally. In this paper, we demonstrate SWG filters in SOI based on BGs and racetrack resonators.

2. SWG BGs

Fig. 1a compares the schematic of an SWG waveguide and an SWG BG in SOI. The effective index of the SWG waveguide depends on the duty cycle \( f = a/\Lambda \), where \( a \) is the width of the high index medium (here Si) and \( \Lambda \) is the period. An SWG BG can then be realized by interleaving two SWG waveguides with different duty cycles \( f_1 = a_1/\Lambda_1 \) and \( f_2 = a_2/\Lambda_2 \); the SWG BG has a period of \( \Lambda_1 + \Lambda_2 \). By varying \( f_1 \) and \( f_2 \), we create a periodic variation in the effective index and can thus obtain Bragg reflection. The cross-section of the SWG waveguides used in our experiments is shown in Fig. 1b. The width of the silicon layer is 500 nm and the thickness is 220 nm; it sits on top of a 3 μm thick buried oxide (BOX) layer and is covered by an index-matched 2 μm thick oxide cladding. We consider SWG waveguides with \( \Lambda = 280 \) nm to obtain a transmission window spanning the C-band and 1,000 periods. Two SWG tapers are used to convert light into (and from) a Bloch mode before (and after) propagating through the SWG waveguide [5]. Vertical grating couplers (VGCs) [6] optimized for TE transmission are used to couple light in and out of the device. A compact Y-branch is used to extract the reflection response of the SWG BGs. The devices were fabricated using electron beam lithography with a single full etch. Figs. 1c and 1d show the SEM image of an SWG BG (before deposition of the top oxide cladding) and the full layout of the device.

![Fig. 1. (a) Top view of an SWG waveguide and an SWG BG, (b) device cross-section, (c) SEM of the fabricated SWG BG prior to top oxide cladding deposition, and (d) full layout.](image)

Fig. 2a shows the measured transmission response of the SWG waveguide without and with an index modulation based on \( f_1 = 50\% \) and \( f_2 = 48\% \). There are no spectral features within the transmission window of the SWG waveguide; on the other hand, the SWG BG exhibits a clear rejection peak at a resonant wavelength of 1546.8 nm. Fig. 2b shows a zoom of the transmission and reflection responses around 1546.8 nm: the transmission loss is -12 dB corresponding to a peak reflectivity of 90.4%; the 3 dB bandwidth is 0.5 nm. The total fiber-to-fiber loss is typically 15 dB (input port to reflection port).
calculated using the total perimeter of the racetrack and effective waveguide index of 2.47. At 1527.9 nm, near critical coupling is achieved and we obtain an extinction ratio (ER) as high as 33 dB; the 3 dB bandwidth is 1 nm. The amplitude and ER variations are due to the wavelength dependence of the directional coupling.

3. SWG Racetrack Resonators

In [5], we demonstrated SWG ring resonators that are under coupled. Racetrack resonators are preferable since the coupling region between the resonator and bus waveguides is longer compared to ring configurations, thereby allowing for greater control over coupling efficiency. Figs. 3a and 3b show the device layout and a zoom of the SWG racetrack resonator. The length of the directional coupling section is \( L \) and the separation between the bus and racetrack waveguides (center-to-center) is denoted \( g \); these two parameters determine the coupling efficiency and operating condition (under coupled, critically coupled, over coupled) of the resonator. The curved section has a radius \( r \) and the total perimeter of the racetrack which determines the free spectral range (FSR) is \( 2\pi + 2L \). We use the same waveguide cross-section and VGC designs as for the SWG BGs. The SWG period is 300 nm and the duty cycle is 50% (to obtain a transmission window in the C-band).

Fig. 4. (a) Measured through and drop responses of a SWG racetrack resonator with \( r = 20 \, \mu m, L = 40 \, \mu m, \) and \( g = 600 \, nm \). (b) Zoom in at one resonance at 1527.85 nm.

4. Conclusions

We have demonstrated and experimentally verified SWG filters in SOI based on BGs and racetrack resonators. The SWG BG exhibits high reflectivity (> 90%) and a relatively narrow reflection bandwidth (0.5 nm). The SWG racetrack resonator exhibits an extinction ratio as high as 33 dB. Based on the SWG resonators, SWG modulators and higher-order filters can be realized. These filters are compatible with other SWG building blocks and can be used to develop more complex structures with enhanced functionality for applications in communications and sensing.

Acknowledgments

The devices were fabricated by R. Bojko at the University of Washington Nanofabrication Facility, a member of the NSF National Nanotechnology Infrastructure Network.

References